

Case proměnná vázba rigózně

$$H \longrightarrow \bar{S}^\dagger H S$$

interakční část $H \longrightarrow \bar{S}^\dagger (H - H_0) S$

$$\frac{\partial}{\partial t} |\Psi(t)\rangle = -\frac{i}{\hbar} \begin{pmatrix} 0 & J \cos \omega t \\ J \cos \omega t & \varepsilon \end{pmatrix} |\Psi(t)\rangle$$

$$H = H_0 + H_I \xrightarrow{\text{I.O.}} H_I^{(I)}(t) = \begin{pmatrix} 0 & J e^{i\omega t} \\ J e^{-i\omega t} & 0 \end{pmatrix}$$

Inverze interakční části

$$|\Psi(t)\rangle = U_0(t, t_0) |\tilde{\Psi}(t)\rangle \quad \text{a nezávisle} \quad H_0 = \begin{pmatrix} 0 & 0 \\ 0 & \frac{1}{2}\hbar\omega \end{pmatrix}$$

$$|\tilde{\Psi}(t)\rangle = U_0(t, t_0) |\Psi(t)\rangle \quad U_0(t, t_0) = \begin{pmatrix} 1 & 0 \\ 0 & e^{-i\omega t} \end{pmatrix}$$

$$\begin{aligned} \frac{\partial}{\partial t} |\tilde{\Psi}(t)\rangle &= \left(\frac{\partial}{\partial t} U_0(t, t_0) \right) |\Psi(t)\rangle + U_0(t, t_0) \frac{\partial}{\partial t} |\Psi(t)\rangle \\ &= -\frac{i}{\hbar} H_0 U_0(t, t_0) |\Psi(t)\rangle + U_0(t, t_0) \left(-\frac{i}{\hbar} H(t) \right) |\Psi(t)\rangle \\ &\quad - \frac{i}{\hbar} U_0(t, t_0) H(t) U_0(t, t_0) |\Psi(t)\rangle \\ &= -\frac{i}{\hbar} H^I(t) |\tilde{\Psi}(t)\rangle \end{aligned}$$

$$H^{(I)}(t) = \begin{pmatrix} 1 & 0 \\ 0 & e^{-i\Omega t} \end{pmatrix} \begin{pmatrix} 0 & J \cos \omega t \\ J \cos \omega t & \varepsilon \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & e^{i\Omega t} \end{pmatrix}$$

$$= \begin{pmatrix} 0 & J \cos \omega t \\ J e^{-i\Omega t} \cos \omega t & \varepsilon e^{-i\Omega t} \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & e^{i\Omega t} \end{pmatrix} =$$

$$= \begin{pmatrix} 0 & J e^{i\Omega t} \cos \omega t \\ J e^{-i\Omega t} \cos \omega t & \varepsilon \end{pmatrix}$$

$$e^{-i\Omega t} \cos \omega t = \frac{1}{2} J \begin{pmatrix} e^{i(\omega - \Omega)t} & e^{i(\omega + \Omega)t} \\ + e & \end{pmatrix}$$

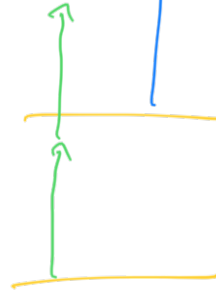
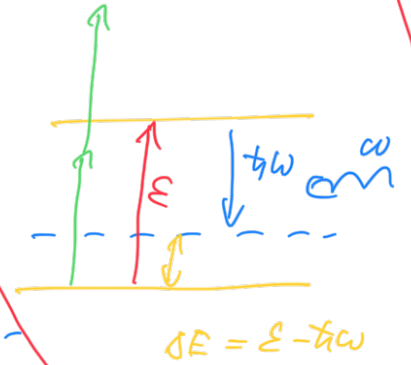
$$H_0 = \begin{pmatrix} 0 & 0 \\ 0 & \varepsilon + \Omega \end{pmatrix}$$

$$H^I(t) = H_0 + H^{(I)}(t) = \begin{pmatrix} 0 & \frac{J}{2} \begin{pmatrix} e^{-i(\omega - \Omega)t} & e^{i(\omega + \Omega)t} \\ + e & \end{pmatrix} \\ \frac{J}{2} \begin{pmatrix} e^{i(\omega - \Omega)t} & e^{-i(\omega + \Omega)t} \\ + e & \end{pmatrix} & \varepsilon + \Omega \end{pmatrix}$$

$$\omega + \Omega = 0 \rightarrow H'(t) \equiv H' = \begin{pmatrix} 0 & \frac{\gamma}{2} \\ \frac{\gamma}{2} & \epsilon - \hbar\omega \end{pmatrix}$$

$$\omega - \Omega = 0$$

$$H'(t) = \begin{pmatrix} 0 & \frac{\gamma}{2} e^{i2\omega t} \\ \frac{\gamma}{2} e^{-i2\omega t} & \epsilon - \hbar\omega \end{pmatrix}$$



Zanedbaním příspěvku oscilujícího se součtovou
frekvencí $\nu = 2\omega$ se říká Aproximace
rotující vlnou — Rotating wave approximation
(RWA)